

## Universal scaling relation in high-temperature superconductors.

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**Recommended and a Commentary by P.W. Anderson, Princeton University.**

The Uemura relation,  $T_c \propto \rho_s$ , although fitting quite well in the underdoped region of high  $T_c$  materials, fails badly for optimal and over doped materials. Though it is often discussed in the context of "ordinary" superconductors, it can hardly hold there since  $\rho_s$  is impurity sensitive and  $T_c$  is not; still less can it possibly hold for c-axis superfluidity, since  $\rho_s(c)$  varies wildly. Where Uemura fits it seems to be compatible with the Lee-Wen idea that superconductivity is destroyed by thermal excitation of quasiparticles.

Homes et al seem to have found a much more nearly universal relation, fitting as well as Uemura where the latter works, and covering also the higher dopings, and, rather amazingly, the c-axis electrodynamics as well as that on the ab plane. They achieve this miracle by dividing the superfluid density by the DC conductivity at the transition temperature  $T_c$ , usually obtained from extrapolation of the infrared conductivity. (This step isn't questionable; direct measurements agree when they exist.) A miracle, to my mind, is that this DC conductivity seems to be constant with doping  $x$  in the underdoped regime, even though  $\rho_s$  is fairly accurately proportional to  $x$ , and even though the resistivity has no simple variation with  $T$ .

The quotient {(superfluid density)/conductivity} defines a relaxation rate  $1/\tau$ , obtained by assuming that the superfluid density contains all the mobile electrons, and that they all relax at the same rate  $1/\tau$ . More formally, this is the result obtained if the conductivity is described by the simplest bubble diagram without vertex corrections, using renormalized single-particle propagators  $G(k, \omega)$  and assuming that the IP of the single-particle self-energy is  $\hbar/\tau$  in the normal state, but the self-energy is purely real in the superconductor. This also amounts to saying that all of the resistivity comes from quasiparticle decay due to strong electron-electron interaction, the familiar basic assumption of either the marginal Fermi-Liquid or the non-FL theories of transport in the normal state--see my book. If this is the case it is interesting, if not original, that the c-axis  $\tau$ , defined this way, should be the same as that defined from the ab plane--as pointed out long ago by N Kumar and collaborator.

What is new, and very germane, is that the scaling law then claims that  $k_B T_c = (\text{const}) x \hbar/\tau$ . If my estimate is right the numerical factor is 1.4 if one uses  $\hbar$ -stroke. This is non-trivial, and must tell us something fairly fundamental about the mechanism for  $T_c$ .

What must be a coincidence gives Homes et al an equivalent, and actually well-known, relationship in the extreme dirty limit of BCS superconductors, though they note that the coefficient is about a factor of 2 different (It's to be found in Tinkham's book, I believe.) But it fails to hold for the clean limit by many orders of magnitude, which must be true since  $T_c$  and  $\rho_s$  are dirt-independent and resistivity not--and of course the mechanism depends not at all on dirt.

There are two accepted kinds of superconducting transition, which until now have been assumed to be the extremes of a one-dimensional manifold where the parameter is the ratio of pair binding energy (the "gap", more or less) to the single particle kinetic energy  $E_f$ . Neither of these obeys anything like Homes' relation (taking the clean limit as more meaningful for the BCS case). For a Bose liquid of bound pairs,  $T_c$  is the degeneracy temperature, proportional to  $n$  in one dimension, and  $\hbar/\tau$  would be  $n$  to the  $3/2$  power for particle collisions,  $n$  to the  $1/2$  for dirt.

I conjecture that the high  $T_c$  transition is of a qualitatively different type and does not belong in the one-dimensional manifold described above. This  $T_c$  occurs when the quasiparticles break apart: it is the deconfinement transition, possibly deconfinement of separate charge and spin excitations; but in any case condensation takes place when the electron as a quasiparticle first begins to have an energy definition comparable to its mean energy  $\Delta E = k_B T$ .